

FOURIER METHOD FOR THE INVERSE COEFFICIENT OF THE PSEUDO-PARABOLIC EQUATION WITH NON-LOCAL BOUNDARY CONDITION

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ABSTRACT. In this work, we have tried to find the inverse coefficient in the quasilinear pseudo-parabolic equation with over determination conditions. It shows the existence, stability of the solution by iteration method and examined numerical solution.

Keywords: Inverse Problem, Quasilinear Parabolic Equation, Crank-Nicolson difference scheme, Overdetermination Data.

AMS Subject Classification: 35K55, 35K70.

1. INTRODUCTION

Recently, there have been a lot of problems with inverse problems that have a lot of applications like chemical diffusion, applications in heat conduction, population dynamics, thermoelasticity, medical science, electrochemistry, engineering, wide scope, chemical engineering. The inverse problem of determining unknown coefficient in a quasi-linear parabolic equation has generated an increasing amount of interest from engineers and scientist [1, 3, 2, 4]. Nonlocal boundary conditions have played a lot of many important roles in heat transfer, thermoelasticity, control theory, mathematical biology, etc.[2, 4]. Let's take the following problem with unknowns (q, u)

$$u_t - u_{xx} - \varepsilon u_{txx} - q(t)u = g(x, t, u), \quad (x, t) \in \Omega, \quad (1)$$

$$u(x, 0) = \theta(x), \quad x \in [0, 1], \quad (2)$$

$$u(0, t) = 0, \quad u_x(0, t) = u_x(1, t), \quad 0 \leq t \leq T, \quad (3)$$

$$h(t) = \int_0^1 u(x, t) dx, \quad 0 \leq t \leq T, \quad (4)$$

Here $\Omega := \{0 < x < 1, 0 < t < T\}$, $\theta(x) \in [0, 1]$ and $g(x, t, u) \in \bar{\Omega} \times (-\infty, \infty)$.

Definition 1.1. $\{q(t), u(x, t)\} \in C[0, T] \times (C^{2,1}(\Omega) \cap C^{1,0}(\bar{\Omega}))$ is called the classical solution.

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2. SOLUTION OF THE INVERSE PROBLEM

Consider the following system of functions on the interval $[0, 1]$: $X_0(x) = x$, $X_{2k-1}(x) = x \cos 2\pi kx$, $X_{2k}(x) = \sin 2\pi kx$, $Y_0(x) = 2$, $Y_{2k-1}(x) = 4 \cos 2\pi kx$, $Y_{2k}(x) = 4(1-x) \sin 2\pi kx$. The systems of these functions arise for the solution of a nonlocal boundary value problem in heat conduction. It is easy to verify that the systems of functions $X_k(x)$ and $Y_k(x)$, $k = 1, 2, 3, \dots$, are biorthonormal on $[0, 1]$. They are also Riesz bases in $L_2[0, 1]$. Let assume the following conditions are ensured. (C1) $h(t) \in C^1[0, T]$, $q(t) \in C[0, T]$. (C2) $\theta(x) \in C^3[0, 1]$, $\theta(x)|_{x=0} = 0$, $\theta_x(x)|_{x=0} = \theta_x(x)|_{x=1}$, (C3) $g(x, t, u)$ is provided following conditions: (1)

$$\left| \frac{\partial^{(n)} g(x, t, u)}{\partial x^n} - \frac{\partial^{(n)} g(x, t, \tilde{u})}{\partial x^n} \right| \leq b(x, t) |u - \tilde{u}|, \quad n = 0, 1, 2,$$

where $b(x, t) \in L_2(\Omega)$, $b(x, t) \geq 0$, (2) $g(x, t, u) \in C^2[0, 1]$, $t \in [0, T]$, (3) $g(x, t, u)|_{x=0} = 0$, $g_x(x, t, u)|_{x=0} = g_x(x, t, u)|_{x=1}$, (4) $g_0(t) \geq 0$, $t \in [0, T]$.

By Fourier method,

$$\begin{aligned} u(x, t) = & 2 \left[\theta_0 e^{-\int_0^t q(\tau) d\tau} + \int_0^t g_0(\tau, u) e^{-\int_\tau^t q(\tau) d\tau} d\tau \right] \\ & + \sum_{k=1}^{\infty} \sin 2\pi kx (\theta_{2k-1} - 4\pi kt \theta_{2k}) e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\int_0^t q(\tau) d\tau} \\ & + \sum_{k=1}^{\infty} x \cos 2\pi kx (\theta_{2k-1} e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\int_0^t q(\tau) d\tau} + \int_0^t g_{2k-1}(\tau, u) e^{\frac{-(2\pi k)^2(t-\tau)}{1+\varepsilon(2\pi k)^2}} e^{-\int_\tau^t q(\tau) d\tau} d\tau) \\ & + \sum_{k=1}^{\infty} \sin 2\pi kx \int_0^t (g_{2k}(\tau, u) - 4\pi k g_{2k-1}(\tau, u)(t-\tau)) e^{\frac{-(2\pi k)^2(t-\tau)}{1+\varepsilon(2\pi k)^2}} e^{-\int_\tau^t q(\tau) d\tau} d\tau. \end{aligned} \quad (5)$$

Under the condition (1)-(3), differentiating (4), we obtain

$$\int_0^1 u_t(x, t) dx = h'(t), \quad 0 \leq t \leq T. \quad (6)$$

(5) and (6) yield

$$q(t) = \frac{\int_0^1 \int_0^1 g(\alpha, \beta, u(\alpha, \beta)) d\alpha d\beta - h'(t)}{h(t)}. \quad (7)$$

Definition 2.1. Let $\{u(t)\} = \{u_0(t), u_{ck}(t), u_{sk}(t), k = 1, \dots, n\}$ is satisfied that

$$\max_{0 \leq t \leq T} \frac{|u_0(t)|}{2} + \sum_{k=1}^{\infty} \left(\max_{0 \leq t \leq T} |u_{ck}(t)| + \max_{0 \leq t \leq T} |u_{sk}(t)| \right) < \infty,$$

by **B**.

$$\|u(t)\| = \max_{0 \leq t \leq T} \frac{|u_0(t)|}{2} + \sum_{k=1}^{\infty} \left(\max_{0 \leq t \leq T} |u_{ck}(t)| + \max_{0 \leq t \leq T} |u_{sk}(t)| \right),$$

be the norm where **B** is Banach space.

Theorem 2.1. *If the assumptions (C1)-(C3) be provided then the problem (1)-(4) has a unique solution.*

Proof. Using iteration to equation (5)

$$u_0^{(N+1)}(t) = u_0^{(0)}(t) + \int_0^t \int_0^1 g(\alpha, \beta, u^{(N)}(\alpha, \beta)) e^{-\int_\beta^t q^{(N)}(\tau) d\tau} d\alpha d\beta, \tag{8}$$

$$u_{2k}^{(N+1)}(t) = u_{2k}^{(0)}(t) + \frac{1}{1 + \varepsilon(2\pi k)^2} \int_0^t \int_0^1 g(\alpha, \beta, u^{(N)}(\alpha, \beta)) \sin 2\pi k\alpha e^{\frac{-(2\pi k)^2(t-\beta)}{1+\varepsilon(2\pi k)^2}} e^{-\int_\beta^t q^{(N)}(\tau) d\tau} d\alpha d\beta,$$

$$u_{2k-1}^{(N+1)}(t) = u_{2k-1}^{(0)}(t) + \frac{1}{1 + \varepsilon(2\pi k)^2} \int_0^t \int_0^1 g(\alpha, \beta, u^{(N)}(\alpha, \beta)) \alpha \cos 2\pi k\alpha e^{\frac{-(2\pi k)^2(t-\beta)}{1+\varepsilon(2\pi k)^2}} e^{-\int_\beta^t q^{(N)}(\tau) d\tau} d\alpha d\beta$$

$$+ \frac{-4\pi k}{1 + \varepsilon(2\pi k)^2} \int_0^t \int_0^1 (t - \beta) g(\alpha, \beta, u^{(N)}(\alpha, \beta)) \sin 2\pi k\alpha e^{\frac{-(2\pi k)^2(t-\beta)}{1+\varepsilon(2\pi k)^2}} e^{-\int_\beta^t q^{(N)}(\tau) d\tau} d\alpha d\beta,$$

$$u_0^{(0)}(t) = \theta_0, u_{2k-1}^{(0)}(t) = (\theta_{2k-1} - 4\pi kt \theta_{2k}) e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\int_0^t q(\tau) d\tau}, u_{2k}^{(0)}(t) = \theta_{2k} e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\int_0^t q(\tau) d\tau}.$$

$$q^{(N)}(t) = \frac{\int_0^t \int_0^1 g(\alpha, \beta, u^{(N)}(\alpha, \beta)) d\alpha d\beta - h'(t)}{h(t)}.$$

From the theorem, we find $u^{(0)}(t) \in \mathbf{B}_1, t \in [0, T]$. For $N = 0$

$$u_0^{(1)}(t) = u_0^{(0)}(t) + \int_0^t \int_0^1 g(\alpha, \beta, u^{(0)}(\alpha, \beta)) e^{-\int_\beta^t q^{(0)}(\tau) d\tau} d\xi d\beta,$$

Adding and subtracting $\int_0^t \int_0^1 f(\alpha, \beta, 0) e^{-\int_\beta^t q^{(0)}(\tau) d\tau} d\alpha d\beta$, we find

$$u_0^{(1)}(t) = \theta_0(t) + \int_0^t \int_0^1 [g(\alpha, \beta, u^{(0)}(\alpha, \beta)) - g(\alpha, \beta, 0)] e^{-\int_\beta^t q^{(0)}(\tau) d\tau} d\alpha d\beta + \int_0^t \int_0^1 g(\alpha, \beta, 0) e^{-\int_\beta^t q^{(0)}(\tau) d\tau} d\alpha d\beta.$$

Applying Cauchy inequality,

$$|u_0^{(1)}(t)| \leq |\theta_0| + \left(\int_0^t d\beta \right)^{\frac{1}{2}} \left(\int_0^t \left\{ \int_0^1 [g(\alpha, \beta, u^{(0)}(\alpha, \beta)) - g(\alpha, \beta, 0)] d\alpha \right\}^2 d\beta \right)^{\frac{1}{2}}$$

$$+ \left(\int_0^t d\beta \right)^{\frac{1}{2}} \left(\int_0^t \left\{ \int_0^1 |g(\alpha, \beta, 0)| d\alpha \right\}^2 d\beta \right)^{\frac{1}{2}},$$

and using Lipschitz condition, we obtain

$$\begin{aligned} \left| u_0^{(1)}(t) \right| &\leq |\theta_0| + \sqrt{t} \left(\int_0^t \left\{ \int_0^1 b(\alpha, \beta) \left| u^{(0)}(\alpha, \beta) \right| d\alpha \right\}^2 d\beta \right)^{\frac{1}{2}} \\ &\quad + \sqrt{t} \left(\int_0^t \left\{ \int_0^1 |g(\alpha, \beta, 0)| d\alpha \right\}^2 d\beta \right)^{\frac{1}{2}}, \end{aligned}$$

and taking maximum, we find:

$$\begin{aligned} \max_{0 \leq t \leq T} \left| u_0^{(1)}(t) \right| &\leq |\theta_0| + \sqrt{T} \|b(x, t)\|_{L_2(\Gamma)} \left\| u^{(0)}(t) \right\|_B \\ &\quad + \sqrt{T} \|g(x, t, 0)\|_{L_2(\Gamma)}. \end{aligned}$$

using the same estimations and Hölder, Bessel inequality and taking maximum,

$$\begin{aligned} \sum_{k=1}^{\infty} \max_{0 \leq t \leq T} \left| u_{2k-1}^{(1)}(t) \right| &\leq \sum_{k=1}^{\infty} |\theta_{2k-1}| + \frac{\sqrt{3}}{12} \|b(x, t)\|_{L_2(D)} \left\| u^{(0)}(t) \right\|_B \\ &\quad + \frac{\sqrt{3}}{12} \|g(x, t, 0)\|_{L_2(\Gamma)}, \end{aligned}$$

and applying the same estimations we obtain,

$$\begin{aligned} \sum_{k=1}^{\infty} \max_{0 \leq t \leq T} \left| u_{2k}^{(1)}(t) \right| &\leq \sum_{k=1}^{\infty} |\theta_{2k-1}| + 4\pi |T| \sum_{k=1}^{\infty} \left| \theta'_{2k} \right| \\ &\quad + \frac{\sqrt{3}}{12} \|b(x, t)\|_{L_2(\Omega)} \left\| u^{(0)}(t) \right\|_B \\ &\quad + \frac{\sqrt{3}}{12} \|g(x, t, 0)\|_{L_2(\Omega)} \\ &\quad + \frac{\sqrt{2}|T|}{4\pi} \|b(x, t)\|_{L_2(\Gamma)} \left\| u^{(0)}(t) \right\|_B \\ &\quad + \frac{\sqrt{2}|T|}{4\pi} \|g(x, t, 0)\|_{L_2(\Omega)}, \end{aligned}$$

and then, we find,

$$\begin{aligned} &\left\| u^{(1)}(t) \right\|_B \\ &\leq 2|\theta_0| + 4 \sum_{k=1}^{\infty} |\theta_{2k-1}| + 4\pi |T| \sum_{k=1}^{\infty} \left| \theta'_{2k} \right| \\ &\quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + 3\sqrt{2}|T|}{3\pi} \right) \|b(x, t)\|_{L_2(\Omega)} \left\| u^{(0)}(t) \right\|_B \\ &\quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + 3\sqrt{2}|T|}{3\pi} \right) \|g(x, t, 0)\|_{L_2(\Omega)}. \end{aligned}$$

$u^{(1)}(t) \in \mathbf{B}$. Same estimations for N ,

$$\begin{aligned} & \left\| u^{(N+1)}(t) \right\|_{\mathbf{B}} \\ & \leq 2|\theta_0| + 8 \sum_{k=1}^{\infty} |\theta_{2k-1}| + 16\pi |T| \sum_{k=1}^{\infty} |\theta'_{2k}| \\ & \quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|b(x, t)\|_{L_2(D)} \left\| u^{(N)}(t) \right\|_{B_1} \\ & \quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|g(x, t, 0)\|_{L_2(D)}. \end{aligned}$$

According to $u^{(N)}(t) \in \mathbf{B}$ and theorem, $u^{(N+1)}(t) \in \mathbf{B}$,

$$\{u(t)\} = \{u_0(t), u_{2k}(t), u_{2k-1}(t), k = 1, 2, \dots\} \in \mathbf{B}.$$

If we used with same estimations, we obtain

$$\begin{aligned} \left\| q^{(N+1)} \right\|_{C[0, T]} & \leq \left| \frac{h'(t)}{h(t)} \right| + \frac{\|b(x, t)\|_{L_2(D)} \left\| u^{(N)}(t) \right\|_B}{|h(t)|} \\ & \quad + \frac{\|g(x, t, 0)\|_{L_2(\Omega)}}{|h(t)|}. \end{aligned}$$

We show that the iterations $u^{(N+1)}(t), q^{(N+1)}$ converge \mathbf{B} and $C[0, T]$, respectively for $N \rightarrow \infty$ Using Cauchy, Bessel, Hölder inequality, Lipschitz condition and taking maximum of both side of the last inequality, we obtain:

$$\begin{aligned} \left\| u^{(1)}(t) - u^{(0)}(t) \right\|_{B_1} & \leq \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|b(x, t)\|_{L_2(\Omega)} \left\| u^{(0)}(t) \right\|_B \\ & \quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|g(x, t, 0)\|_{L_2(\Omega)}. \end{aligned}$$

$$\begin{aligned} A & = \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|b(x, t)\|_{L_2(\Omega)} \left\| u^{(0)}(t) \right\|_B \\ & \quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|g(x, t, 0)\|_{L_2(\Omega)}. \end{aligned}$$

$$\left\| q^{(1)}(t) - q^{(0)}(t) \right\|_{C[0, T]} \leq \frac{\sqrt{T}}{|h(t)|} \left\| u^{(1)}(t) - u^{(0)}(t) \right\|_B \|b(x, t)\|_{L_2(\Omega)},$$

$$\begin{aligned} \left\| u^{(2)}(t) - u^{(1)}(t) \right\|_{B_1} & \leq \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \|b(x, t)\|_{L_2(\Omega)} \left\| u^{(1)}(t) - u^{(0)}(t) \right\|_B \\ & \quad + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \frac{M\sqrt{T}}{|h(t)|} \left\| u^{(1)}(t) - u^{(0)}(t) \right\|_B \|b(x, t)\|_{L_2(\Omega)}. \end{aligned}$$

$$\left\| u^{(2)}(t) - u^{(1)}(t) \right\|_{B_1} \leq \left\{ \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \left(1 + \frac{M\sqrt{T}}{|h(t)|} \right) \right\} A \|b(x, t)\|_{L_2(\Omega)},$$

$$\begin{aligned} \|q^{(2)}(t) - q^{(0)}(t)\|_{C[0,T]} &\leq \frac{\sqrt{T}}{|h(t)|} \|u^{(2)}(t) - u^{(0)}(t)\|_B \|b(x,t)\|_{L_2(\Omega)}, \\ \|u^{(3)}(t) - u^{(2)}(t)\|_{B_1} &\leq \left\{ \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi)|T|}{3\pi} \right) \left(1 + \frac{M\sqrt{T}}{|h(t)|} \right) \right\}^2 \frac{A}{\sqrt{2!}} \|b(x,t)\|_{L_2(\Omega)}. \end{aligned}$$

For N :

$$\begin{aligned} \|q^{(N+1)}(t) - q^{(N)}(t)\|_{C[0,T]} &\leq \frac{\sqrt{T}}{|h(t)|} \|u^{(N+1)}(t) - u^{(N)}(t)\|_B \|b(x,t)\|_{L_2(\Omega)}, \\ \|u^{(N+1)}(t) - u^{(N)}(t)\|_B &\leq \frac{A}{\sqrt{N!}} \left\{ \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + 3\sqrt{2}|T|}{3\pi} \right) \left(1 + \frac{M\sqrt{T}}{|h(t)|} \right) \right\}^N \|b(x,t)\|_{L_2(\Omega)}. \end{aligned}$$

For $N \rightarrow \infty$, $u^{(N+1)}(t)$, $q^{(N+1)}$ are converged. Let show that there exists u and q such that

$$\lim_{N \rightarrow \infty} u^{(N+1)}(t) = u(t), \quad \lim_{N \rightarrow \infty} q^{(N+1)}(t) = q(t).$$

Using same inequality and Gronwall's inequality, we obtain

$$\begin{aligned} \|u(t) - u^{(N+1)}(t)\|_{\mathbf{B}}^2 &\leq \\ &2 \left[\frac{A}{\sqrt{N!}} \left\{ \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi)|T|}{3\pi} \right) \right\}^N \|b(x,t)\|_{L_2(\Omega)}^N \right]^2 \\ &\times \exp 2 \left\{ \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi)|T|}{3\pi} \right) \left(1 + \frac{M\sqrt{T}}{|h(t)|} \right) \right\}^2 \|b(x,t)\|_{L_2(\Omega)}^2. \end{aligned} \quad (9)$$

$$\|q(t) - q^{(N+1)}(t)\|_{C[0,T]} \leq \frac{\sqrt{T}}{|h(t)|} \|u(t) - u^{(N+1)}(t)\|_B \|b(x,t)\|_{L_2(\Omega)},$$

we obtain $u^{(N+1)} \rightarrow u$, $q^{(N+1)} \rightarrow q$, $N \rightarrow \infty$. For the uniqueness, let (u, q) , (v, r) are two solution of (1)-(4). After applying Cauchy, Bessel, Lipschitz, Hölder inequality to $|u(t) - v(t)|$ and $|r(t) - q(t)|$, we obtain

$$\|u(t) - v(t)\|_{\mathbf{B}} \leq \left[\left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi)|T|}{3\pi} \right) \left(1 + \frac{M\sqrt{T}}{|h(t)|} \right) \right] \left(\int_0^t \int_0^1 b^2(\alpha, \beta) |u(\beta) - v(\beta)|^2 d\alpha d\beta \right) \quad (10)$$

$u(t) = v(t)$ and then $r(t) = q(t)$. □

The proof is over.

3. STABILITY OF PROBLEM

Theorem 3.1. *Assumption (C1)-(C3) the solution (q, u) of the problem (1)-(4) depends continuously upon the data θ, h .*

Proof. Let $\Phi = \{\theta, h, g\}$ and $\bar{\Phi} = \{\bar{\theta}, \bar{h}, \bar{g}\}$ be two sets of the data, which satisfy the assumptions (1) – (3). Let us denote $\|\Phi\| = (\|h\|_{C^1[0,T]} + \|\theta\|_{C^3[0,1]} + \|g\|_{C^{3,0}(\bar{\Gamma})})$. Let (q, u) and (\bar{q}, \bar{u}) be solutions of problems (1)-(4).

$$\begin{aligned}
 u - \bar{u} = & 2(\theta_0 - \bar{\theta}_0) e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} + 2\theta_0 \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) \\
 & + 4 \sum_{k=1}^{\infty} (\theta_{2k-1} - \bar{\theta}_{2k-1}) x \cos 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\tau} \int_0^t \bar{q}(\tau) d\tau} \\
 & + 4 \sum_{k=1}^{\infty} \theta_{2k-1} x \cos 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) \\
 & + 4 \sum_{k=1}^{\infty} (\theta_{2k} - \bar{\theta}_{2k}) \sin 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\tau} \int_0^t \bar{q}(\tau) d\tau} \\
 & + 4 \sum_{k=1}^{\infty} \theta_{2k} \sin 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) \\
 & - 16\pi \sum_{k=1}^{\infty} (\theta_{2k-1} - \bar{\theta}_{2k-1}) k t \sin 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\tau} \int_0^t \bar{q}(\tau) d\tau} \\
 & - 16\pi \sum_{k=1}^{\infty} \theta_{2k-1} k t \sin 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) \\
 & + 2 \int_0^t \int_0^1 [g(\alpha, \beta, u(\alpha, \beta)) - g(\alpha, \beta, \bar{u}(\alpha, \beta))] e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\tau} \int_0^t \bar{q}(\tau) d\tau} \\
 & + 2 \int_0^t \int_0^1 g(\alpha, \beta, u(\alpha, \beta)) e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) \\
 & + 4 \sum_{k=1}^{\infty} \int_0^t \int_0^1 [g(\alpha, \beta, u(\alpha, \beta)) - g(\alpha, \beta, \bar{u}(\alpha, \beta))] (1-x) \sin 2\pi k x e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \\
 & - 16\pi \sum_{k=1}^{\infty} [g(\alpha, \beta, u(\alpha, \beta)) - g(\alpha, \beta, \bar{u}(\alpha, \beta))] e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} e^{-\frac{t}{\tau} \int_0^t \bar{q}(\tau) d\tau} (t-\tau) \cos 2\pi k x \\
 & - 16\pi \sum_{k=1}^{\infty} [g(\alpha, \beta, u(\alpha, \beta)) - g(\alpha, \beta, \bar{u}(\alpha, \beta))] e^{\frac{-(2\pi k)^2 t}{1+\varepsilon(2\pi k)^2}} \left(e^{-\frac{t}{\beta} \int_0^t q(\tau) d\tau} - e^{-\frac{t}{\beta} \int_0^t \bar{q}(\tau) d\tau} \right) (t-\tau) \cos 2\pi k x.
 \end{aligned}$$

By using same estimations, we obtain:

$$\begin{aligned}
 \|u - \bar{u}\|_B \leq & M_3 \|\theta - \bar{\theta}\| \\
 & + \left(\frac{2\pi\sqrt{3} + 6\pi\sqrt{T} + (1 + 2\sqrt{2}\pi) |T|}{3\pi} \right) \left(\int_0^t \int_0^1 b^2(\alpha, \beta) |u(\beta) - \bar{u}(\beta)|^2 d\alpha d\beta \right)^{\frac{1}{2}}.
 \end{aligned} \tag{11}$$

$$\|q - \bar{q}\|_B \leq M_1 \left\| h(t) - \overline{h(t)} \right\|_{C^1[0,T]} + M_2 \left\| u(t) - \overline{u(t)} \right\|_B \|b(x, t)\|_{L_2(\Omega)},$$

applying Gronwall's inequality, we obtain:

$$\begin{aligned} \|u - \bar{u}\|_B &\leq 2M_3 \|\Phi - \bar{\Phi}\|^2 \\ &\quad \times \exp 2 \left(\int_0^t \int_0^1 b^2(\alpha, \beta) d\alpha d\beta \right). \end{aligned}$$

For $\Phi \rightarrow \bar{\Phi}$ then $u \rightarrow \bar{u}$. Hence $q \rightarrow \bar{q}$. □

4. NUMERICAL PROCEDURE FOR THE NONLINEAR PROBLEM (1)-(4)

An iteration algorithm for the linearization of the problem (1)-(4):

$$\frac{\partial u^{(n)}}{\partial t} - \frac{\partial^2 u^{(n)}}{\partial x^2} - \varepsilon \frac{\partial^3 u^{(n)}}{\partial t \partial x^2} - q(t)u^{(n)} = g(x, t, u^{(n-1)}) \quad (x, t) \in D. \quad (12)$$

$$u^{(n)}(x, 0) = \theta(x), \quad x \in [0, 1]. \quad (13)$$

$$u^{(n)}(0, t) = 0, \quad t \in [0, T]. \quad (14)$$

$$u_x^{(n)}(0, t) = u_x^{(n)}(1, t), \quad t \in [0, T]. \quad (15)$$

Let $u^{(n)}(x, t) = v(x, t)$ and $g(x, t, u^{(n-1)}) = \tilde{g}(x, t)$ then a new linear problem :

$$\frac{\partial v}{\partial t} - \frac{\partial^2 v}{\partial x^2} - \varepsilon \frac{\partial^3 v}{\partial t \partial x^2} - q(t)v = \tilde{g}(x, t) \quad (x, t) \in D. \quad (16)$$

$$v(0, t) = 0, \quad t \in [0, T]. \quad (17)$$

$$v_x(0, t) = v_x(1, t), \quad t \in [0, T]. \quad (18)$$

$$v(x, 0) = \theta(x), \quad x \in [0, 1]. \quad (19)$$

we use the method of the linearization and the finite difference method to solve (16)-(19).

We subdivide the intervals $[0, 1]$ and $[0, T]$ into subintervals N_x and N_t of equal lengths $h = \frac{1}{N_x}$ and $\tau = \frac{T}{N_t}$, respectively. We use the Crank-Nicolson scheme which is absolutely stable and has a second-order accuracy in h and a first-order accuracy in τ . The Crank-Nicolson scheme for (16)-(19) is as follows:

$$\begin{aligned} \frac{1}{\tau} \left(v_i^{j+1} - v_i^j \right) &= \frac{1}{h^2} \left(v_{i-1}^{j+1} - 2v_i^{j+1} + v_{i+1}^{j+1} \right) \\ &\quad + \varepsilon \frac{1}{2h^2\tau} \left[\left(v_{i-1}^{j+1} - 2v_i^{j+1} + v_{i+1}^{j+1} \right) - \left(v_{i-1}^j - 2v_i^j + v_{i+1}^j \right) \right] \\ &\quad + q^{j+1}v_i^{j+1} + \tilde{g}_i^{j+1}, \end{aligned} \quad (20)$$

$$v_i^0 = \theta_i, \quad (21)$$

$$v_0^j = 0, \quad (22)$$

$$v_{N_x+1}^j = v_1^j + v_{N_x}^j, \quad (23)$$

where $1 \leq i \leq N_x$ and $0 \leq j \leq N_t$ are the indices for the spatial and time steps respectively, $v_i^j = v(x_i, t_j)$, $\theta_i = \theta(x_i)$, $q^j = q(t_j)$, $\tilde{g}_i^j = \tilde{g}(x_i, t_j)$, $x_i = ih$, $t_j = j\tau$.

$$q^j = \frac{((h^{j+1} - h^j) / \tau) - \tilde{g}_i^{j+1}}{h_i^j}, \quad (24)$$

where $h_i^j = h(t_j)$, $j = 0, 1, \dots, N_t$.

The system can be solved by the Gauss elimination method, and v_i^j is determined.

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