# NONHOLONOMIC FRAMES FOR FINSLER SPACE WITH DEFORMED MATSUMOTO METRIC 

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Abstract. The purpose of present paper is to find the nonholonomic frames for the deformed Matsumoto type metric which are given in the forms
I. $\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \alpha=\frac{\alpha^{3}}{\alpha-\beta}$
II. $\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \beta=\frac{\alpha^{2} \beta}{\alpha-\beta}$
where $\alpha^{2}=a_{i j}(x) y^{i} y^{j}$ and $\beta=b_{i}(x) y^{i}$. The first metric of the above deformation is obtained by the product of Matsumoto and Riemannian metric and second one is the product of Matsumoto and 1-form metric.

Keywords: Riemannian metric, one form metric, Matsumoto metric, GL-metric, nonholonomic Finsler frame.

AMS Subject Classification: 53C60

## 1. Introduction

P.R. Holland [1] [2] studies about the nonholonomic frame on space time which was based on the consideration of a charged particle moving in an external electromagnetic field in the year 1982. In the year 1987, R.S. Ingarden [3] was the first person, to point out that the Lorentz force law can be written in above case as geodesic equation on a Finsler space called Randers space. Further in 1995, R.G. Beil [5] [6] have studied a gauge transformation viewed as a nonholonomic frame on the tangent bundle of a four dimensional base manifold. The geometry that follows from these considerations gives a unified approach to gravitation and gauge symmetries.

In the present paper, we have used the common Finsler idea to study the existence of a nonholonomic frame on the vertical sub bundle VTM of the tangent bundle of a base manifold M. In this case we have considered the fundamental tensor field might be the deformation of two different special Finsler spaces from the $(\alpha, \beta)$ - metrics. First we consider a nonholonomic frame for a Finsler space with $(\alpha, \beta)$ - metrics such as first product of Matsumoto metric[11] and Riemannian metric and second is the product of Matsumoto metric[11] and 1-form metric. Further, we obtain a corresponding frame for each of these

[^0]two Finsler deformations. This is an extension work of Ioan Bucataru and Radu Miron [10], Tripathi $[14,16]$ and Narasimhamurthy [15].

## 2. Preliminaries

An important class of Finsler spaces is the class of Finsler spaces with $(\alpha, \beta)$-metrics [11].
Definition 2.1. A Finsler space $F^{n}=\{M, F(x, y)\}$ is called with ( $\alpha, \beta$ )-metric if there exists a 2-homogeneous function $L$ of two variables such that the Finsler metric $F: T M \rightarrow$ $R$ is given by

$$
\begin{equation*}
F^{2}(x, y)=L\{\alpha(x, y), \beta(x, y)\} \tag{1}
\end{equation*}
$$

where $\alpha^{2}(x, y)=a_{i j}(x) y^{i} y^{j}, \alpha$ is a Riemannian metric on the manifold $M$, and $\beta(x, y)=$ $b_{i}(x) y^{i}$ is a 1 -form on $M$.

The first Finsler spaces with $(\alpha, \beta)$-metrics were introduced by the physicist G. Randers in 1940, are called Randers spaces [4]. Recently, R.G. Beil suggested a more general case by considering, $a_{i j}(x)$ the components of a Riemannian metric on the base manifold M, $a(x, y)>0$ and $b(x, y) \geq 0$ Two functions on TM and $B(x, y)=B_{i}(x, y)\left(d x^{i}\right)$ a vertical 1-form on TM. Then

$$
g_{i j}(x, y)=a(x, y) a_{i j}(x)+b(x, y) B_{i}(x, y) B_{j}(x, y)
$$

Now a days the above generalized Lagrange metric is known as the Beil metric. The metric tensor $g_{i j}$ is also known as a Beil deformation of the Riemannian metric $a_{i j}$. It has been studied and applied by R. Miron and R.K. Tavakol in General Relativity for $a(x, y)=\exp (2 \sigma(x, y))$ and $\mathrm{b}=0$. The case $\mathrm{a}(\mathrm{x}, \mathrm{y})=1$ with various choices of b and $B_{i}$ was introduced and studied by R.G. Beil for constructing a new unified field theory [6]. Further Bucataru [12] considered the class of Lagrange spaces with ( $\alpha, \beta$ )-metric and obtained some new and interesting results in the year 2002.

A unified formalism which uses a nonholonomic frame on space time, a sort of plastic deformation, arising from consideration of a charged particle moving in an external electromagnetic field in the background space time viewed as a strained mechanism studied by P. R. Holland. If we do not ask for the function $L$ to be homogeneous of order two with respect to the $(\alpha, \beta)$ variables, then we have a Lagrange space with $(\alpha, \beta)$-metric. Next we defined some different Finsler space with $(\alpha, \beta)$-metrics.

Further consider $g_{i j}=\frac{1}{2} \frac{\partial^{2} F^{2}}{\partial y^{2} \partial y^{j}}$ the fundamental tensor of the Randers space(M,F). Taking into account the homogeneity of a and F we have the following formulae:

$$
\begin{gather*}
p^{i}=\frac{1}{a} y^{i}=a^{i j} \frac{\partial \alpha}{\partial y^{j}} ; \quad p_{i}=a_{i j} p^{j}=\frac{\partial \alpha}{\partial y^{i}} \\
l^{i}=\frac{1}{L} y^{i}=g^{i j} \frac{\partial l}{\partial y^{i}} ; l_{i}=g_{i j} l^{j}=\frac{\partial L}{\partial y^{i}}=P_{i}+b_{i}  \tag{2}\\
l^{i}=\frac{1}{L} p^{i} ; l^{i} l_{i}=p^{i} p_{i}=1 ; l^{i} p_{i}=\frac{\alpha}{L} ; p^{i} l_{i}=\frac{L}{\alpha} \\
b_{i} P^{i}=\frac{\beta}{\alpha} ; b_{i} l^{i}=\frac{\beta}{L}
\end{gather*}
$$

with respect to these notations, the metric tensors $a_{i j}$ and $g_{i j}$ are related by [13],

$$
\begin{equation*}
g_{i j}(x, y)=\frac{L}{\alpha} a_{i j}+b_{i} P_{j}+P_{i} b_{j}-\frac{\beta}{\alpha} p_{i} p_{j}=\frac{L}{\alpha}\left(a_{i j}-p_{i} p_{j}\right)+l_{i} l_{j} \tag{3}
\end{equation*}
$$

Theorem 2.1. [10]: For a Finsler space $(M, F)$ consider the metric with the entries:

$$
\begin{equation*}
Y_{j}^{i}=\sqrt{\frac{\alpha}{L}}\left(\delta_{j}^{i}-l^{i} l_{j}+\sqrt{\frac{\alpha}{L}} p^{i} p_{j}\right) \tag{4}
\end{equation*}
$$

defined on $T M$. Then $Y_{j}=Y_{j}^{i}\left(\frac{\partial}{\partial y^{i}}\right), \quad j \in 1,2,3, \ldots, n$ is a non holonomic frame.
Theorem 2.2. [7]: With respect to frame the holonomic components of the Finsler metric tensor $a_{\alpha \beta}$ is the Randers metric $g_{i j}$, i.e,

$$
\begin{equation*}
g_{i j}=Y_{i}^{\alpha} Y_{j}^{\beta} a_{\alpha \beta} \tag{5}
\end{equation*}
$$

Throughout this section we shall rise and lower indices only with the Riemannian metric $a_{i j}(x)$ that is $y_{i}=a_{i j} y^{j}, \beta^{i}=a^{i j} b_{j}$, and so on. For a Finsler space with $(\alpha, \beta)$-metric $F^{2}(x, y)=L\{\alpha(x, y), \beta(x, y)\}$ we have the Finsler invariants [13].

$$
\begin{equation*}
\rho_{1}=\frac{1}{2 \alpha} \frac{\partial L}{\partial \alpha} ; \rho_{0}=\frac{1}{2} \frac{\partial^{2} L}{\partial \beta^{2}} ; \rho_{-1}=\frac{1}{2 \alpha} \frac{\partial^{2} L}{\partial \alpha \partial \beta} ; \rho_{-2}=\frac{1}{2 \alpha^{2}}\left(\frac{\partial^{2} L}{\partial \alpha^{2}}-\frac{1}{\alpha} \frac{\partial L}{\partial \alpha}\right) \tag{6}
\end{equation*}
$$

where subscripts $1,0,-1,-2$ gives us the degree of homogeneity of these invariants. For a Finsler space with $(\alpha, \beta)$-metric we have,

$$
\begin{equation*}
\rho_{-1} \beta+\rho_{-2} \alpha^{2}=0 \tag{7}
\end{equation*}
$$

with respect to the notations we have that the metric tensor $g_{i j}$ of a Finsler space with $(\alpha, \beta)$-metric is given by [13].

$$
\begin{equation*}
g_{i j}(x, y)=\rho a_{i j}(x)+\rho_{0} b_{i}(x)+\rho_{-1}\left\{b_{i}(x) y_{j}+b_{j}(x) y_{i}\right\}+\rho_{-2} y_{i} y_{j} \tag{8}
\end{equation*}
$$

From (8) we can see that $g_{i j}$ is the result of two Finsler deformations:

$$
\begin{array}{ll}
\text { I. } & a_{i j} \rightarrow h_{i j}=\rho a_{i j}+\frac{1}{\rho_{-2}}\left(\rho_{-1} b_{i}+\rho_{-2} y_{i}\right)\left(\rho_{-1} b_{j}+\rho_{-2} y_{j}\right) \\
\text { II. } & h_{i j} \rightarrow g_{i j}=h_{i j}+\frac{1}{\rho_{-2}}\left(\rho_{0} \rho_{-1}-\rho_{-1}^{2}\right) b_{i} b_{j} \tag{9}
\end{array}
$$

The nonholonomic Finsler frame that corresponding to the $I^{\text {st }}$ deformation (9) is according to the theorem (7.9.1) in [10], given by,

$$
\begin{equation*}
X_{j}^{i}=\sqrt{\rho} \delta_{j}^{i}-\frac{1}{\beta^{2}}\left\{\sqrt{\rho}+\sqrt{\rho+\frac{\beta^{2}}{\rho_{-2}}}\right\}\left(\rho_{-1} b^{i}+\rho_{-2} y^{i}\right)\left(\rho_{-1} b_{j}+\rho_{-2} y_{j}\right) \tag{10}
\end{equation*}
$$

where $B^{2}=a_{i j}\left(\rho_{-1} b^{i}+\rho_{-2} y^{i}\right)\left(\rho_{-1} b_{j}+\rho_{-2} y_{j}\right)=\rho_{-1}^{2} b^{2}+\beta \rho_{-1} \rho_{-2}$.
This metric tensor $a_{i j}$ and $h_{i j}$ are related by,

$$
\begin{equation*}
h_{i j}=X_{i}^{k} X_{j}^{l} a_{k l} \tag{11}
\end{equation*}
$$

Again the frame that corresponds to the $I I_{n d}$ deformation (9) given by,

$$
\begin{equation*}
Y_{j}^{i}=\delta_{j}^{i}-\frac{1}{C^{2}}\left\{1 \pm \sqrt{1+\left(\frac{\rho_{-2} C^{2}}{\rho_{0} \rho_{-2}-\rho_{-1}^{2}}\right)}\right\} b^{i} b_{j} \tag{12}
\end{equation*}
$$

where $C^{2}=h_{i j} b^{i} b^{j}=\rho b^{2}+\frac{1}{\rho_{-2}}\left(\rho_{-1} b^{2}+\rho_{-2} \beta\right)^{2}$.
The metric tensor $h_{i j}$ and $g_{i j}$ are related by the formula;

$$
\begin{equation*}
g_{m n}=Y_{m}^{i} Y_{n}^{j} h_{i j} \tag{13}
\end{equation*}
$$

Theorem 2.3. [10]: $\operatorname{Let} F^{2}(x, y)=L\{\alpha(x, y), \beta(x, y)\}$ be the metric function of a Finsler space with $(\alpha, \beta)$ metric for which the condition (7) is true. Then

$$
V_{j}^{i}=X_{k}^{i} Y_{j}^{k}
$$

is a nonholonomic Finsler frame with $X_{k}^{i}$ and $Y_{j}^{k}$ are given by (10) and (12) respectively.

## 3. Nonholonomic frames for Finsler space with deformed Matsumoto Metric

In this section we consider two cases of nonholonomic Finlser frames with special $(\alpha, \beta)$ metrics, such a $I^{s t}$ Finsler frame product of Matusmoto metric and Riemannian metric and $I I^{n d}$ Finsler frame product of Matsumoto metric and 1-form metric.
3.1. Nonholonomic frame for $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \alpha=\frac{\alpha^{3}}{\alpha-\beta}$. In the first case, for a Finsler space with the fundamental function $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \alpha=\frac{\alpha^{3}}{\alpha-\beta}$ the Finsler invariants (6) are given by

$$
\begin{align*}
& \rho_{1}=\frac{2 \alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{2}}, \quad \rho_{0}=\frac{\alpha^{3}}{(\alpha-\beta)^{3}}  \tag{14}\\
& \rho_{-1}=\frac{\alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{3}}, \quad \rho_{-2}=\frac{3 \beta^{2}-\alpha \beta}{2 \alpha(\alpha-\beta)^{3}} \\
& B^{2}=\frac{(\alpha-3 \beta)^{2}\left(\alpha^{2} b^{2}-\beta^{2}\right)}{4(\alpha-\beta)^{6}}
\end{align*}
$$

Using (14) in (10) we have,

$$
\begin{align*}
& X_{j}^{i}=\sqrt{\frac{2 \alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{2}}} \delta_{j}^{i}-\frac{\left(\alpha^{2}-3 \alpha \beta\right)^{2}}{4 \beta^{2}(\alpha-\beta)^{6}}\left[\sqrt{\frac{2 \alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{2}}}+\right.  \tag{15}\\
& \left.\quad \sqrt{\frac{2 \alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{2}}+\frac{2 \alpha \beta(\alpha-\beta)^{3}}{(3 \beta-\alpha)}}\right]\left(b^{i}-\frac{\beta}{\alpha^{2}} y^{i}\right)\left(b_{j}-\frac{\beta}{\alpha^{2}} y_{j}\right)
\end{align*}
$$

Again using (14) in (12) we have,

$$
\begin{equation*}
Y_{j}^{i}=\delta_{j}^{i}-\frac{1}{C^{2}}\left\{1 \pm \sqrt{1+\frac{2 \beta(\alpha-\beta)^{2} C^{2}}{\alpha^{3}}}\right\} b^{i} b_{j} \tag{16}
\end{equation*}
$$

where $C^{2}=\frac{2 \alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{2}} b^{2}+\frac{(3 \beta-\alpha)}{2 \alpha \beta(\alpha-\beta)^{3}}\left(\alpha^{2} b^{2}-\beta^{2}\right)^{2}$
Theorem 3.1. Let $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \alpha=\frac{\alpha^{3}}{\alpha-\beta}$ be the metric function of a Finsler space with ( $\alpha$, $\beta$ ) metric for which the condition (7) is true. Then

$$
V_{j}^{i}=X_{k}^{i} Y_{j}^{k}
$$

is nonholonomic Finsler Frame with $X_{k}^{i}$ and $Y_{j}^{k}$ are given by (15) and (16) respectively.
3.2. Nonholonomic frame for $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \beta=\frac{\alpha^{2} \beta}{\alpha-\beta}$. In the second case, for a Finsler space with the fundamental function $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \beta=\frac{\alpha^{2} \beta}{\alpha-\beta}$ the Finsler invariants (6) are
given by

$$
\begin{array}{r}
\rho_{1}=\frac{\alpha \beta-2 \beta^{2}}{2(\alpha-\beta)^{2}}, \quad \rho_{0}=\frac{\alpha^{3}}{(\alpha-\beta)^{3}}  \tag{17}\\
\rho_{-1}=\frac{\alpha^{2}-3 \alpha \beta}{2(\alpha-\beta)^{3}}, \quad \rho_{-2}=\frac{3 \beta^{2}-\alpha \beta}{2 \alpha(\alpha-\beta)^{3}} \\
B^{2}=\frac{(\alpha-3 \beta)^{2}\left(\alpha^{2} b^{2}-\beta^{2}\right)}{4(\alpha-\beta)^{6}}
\end{array}
$$

Using (17) in (10) we have,

$$
\begin{align*}
& X_{j}^{i}=\sqrt{\frac{\alpha \beta-2 \beta^{2}}{2(\alpha-\beta)^{2}}} \delta_{j}^{i}-\frac{\left(\alpha^{2}-3 \alpha \beta\right)^{2}}{4 \beta^{2}(\alpha-\beta)^{6}}\left[\sqrt{\frac{\alpha \beta-2 \beta^{2}}{2(\alpha-\beta)^{2}}}+\right.  \tag{18}\\
& \left.\sqrt{\frac{\alpha \beta-2 \beta^{2}}{2(\alpha-\beta)^{2}}+\frac{2 \alpha \beta(\alpha-\beta)^{3}}{(3 \beta-\alpha)}}\right]\left(b^{i}-\frac{\beta}{\alpha^{2}} y^{i}\right)\left(b_{j}-\frac{\beta}{\alpha^{2}} y_{j}\right)
\end{align*}
$$

Again using (17) in (12) we have,

$$
\begin{equation*}
Y_{j}^{i}=\delta_{j}^{i}-\frac{1}{C^{2}}\left\{1 \pm \sqrt{1+\frac{2 \beta(\alpha-\beta)^{2} C^{2}}{\alpha^{3}}}\right\} b^{i} b_{j} \tag{19}
\end{equation*}
$$

where $C^{2}=\frac{\alpha \beta-2 \beta^{2}}{2(\alpha-\beta)^{2}} b^{2}+\frac{(3 \beta-\alpha)}{2 \alpha \beta(\alpha-\beta)^{3}}\left(\alpha^{2} b^{2}-\beta^{2}\right)^{2}$
Theorem 3.2. Let $L=\left(\frac{\alpha^{2}}{\alpha-\beta}\right) \beta=\frac{\alpha^{2} \beta}{\alpha-\beta}$ be the metric function of a Finsler space with $(\alpha, \beta)$ metric for which the condition (7) is true. Then

$$
V_{j}^{i}=X_{k}^{i} Y_{j}^{k}
$$

is nonholonomic Finsler Frame with $X_{k}^{i}$ and $Y_{j}^{k}$ are given by (18) and (19) respectively.

## 4. Conclusions

Nonholonomic frame relates a semi-Riemannian metric (the Minkowski or the Lorentz metric) with an induced Finsler metric. Antonelli and Bucataru ([7][8]), has been determined such a nonholonomic frame for two important classes of Finsler spaces that are dual in the sense of Randers and Kropina spaces [9]. As Randers and Kropina spaces are members of a bigger class of Finsler spaces, namely the Finsler spaces with $(\alpha, \beta)$-metric, it appears a natural question: Does how many Finsler space with $(\alpha, \beta)$-metrics have such a nonholonomic frame? The answer is yes, there are many Finsler space with $(\alpha, \beta)$-metrics.

In this work, we consider the Matsumoto Finsler metrics, Riemannian metric and 1form metric we determine the nonholonomic Finsler frames. But, in Finsler geometry, there are many $((\alpha, \beta)$-metrics, in future work we can determine the frames for them also.

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